Part 2a

The total energy radiated per second = $4\pi R^2 \sigma T^4$, where σ is the Stephan-Boltzmann constant. The energy incident on a unit area on earth per second is;

P =
$$\frac{4\pi R^2 \sigma T^4}{4\pi \ell^2}$$
 yielding, $R = \left(P/\sigma T^4\right)^{1/2} \ell$ (1) (0.8 pts)

The energy of a photon is $hf=hc/\lambda$. The equivalent mass of a photon is $h/c\lambda$. Conservation of photon energy:

$$\frac{hc}{\lambda_0} - \frac{Gm_0}{R} \cdot \frac{h}{c\lambda_0} = \frac{hc}{\lambda}$$
 (0.8 pts)

yielding

$$R = \frac{Gm_0(\lambda_0 + \Delta\lambda)}{c^2 \Lambda\lambda}$$
 (2)

and (2) yields,

$$m_0 = \frac{c^2 \Delta \lambda \left(P / \sigma T^4 \right)^{1/2}}{G(\lambda_0 + \Delta \lambda)} \ell \quad (3)$$

The stars are rotating around the center of mass with equal angular speeds:

$$\omega = (2\pi/2\tau) = \pi/\tau$$
 (4) (0.2 pts)

The equilibrium conditions for the stars are;

$$\frac{GMm_0}{(r_1 + r_2)^2} = m_0 r_1 \omega^2 = Mr_2 \omega^2$$
 (5)

with

$$r_1 = \ell \frac{\Delta \theta}{2}$$
, $r_2 = \ell \frac{\Delta \phi}{2}$ (6)

Substituting (3), (4) and (6) into (5) yields

$$\ell = \left(\frac{8c^2\Delta\lambda \left(P/\sigma T^4\right)^{1/2}}{\Delta\phi(\pi/\tau)^2 \left(\lambda_0 + \Delta\lambda\right) \left(\Delta\theta + \Delta\phi\right)^2}\right)^{1/2}.$$
 (0.8 pts)

Part 2b

Conservation of angular momentum for the ordinary star;

$$mr^2\omega = m_0 r_0^2 \omega_0$$
 (7) (0.6 pts.)

Conservation of angular momentum for dm:

$$r^2\omega dm = r_f^2\omega_f dm$$
 (8)

where ω_{f} is the angular velocity of the ring. Equilibrium in the original state yields,

$$\omega_0 = \left(\frac{GM}{r_0^3}\right)^{1/2}$$
 (9)

and (7), (8) and (9) give,

$$\omega = \frac{m_0 r_0}{mr^2} \left(\frac{GM}{r_0}\right)^{1/2}, \ \omega_f = \frac{m_0 r_0}{mr_f^2} \left(\frac{GM}{r_0}\right)^{1/2}$$
 (10)

Conservation of energy for dm;

$$\frac{1}{2}dm\left(v_0^2 + r^2\omega^2\right) - \frac{GM \text{ dm}}{r} = \frac{1}{2}dm r_f^2 \omega_f^2 - \frac{GM \text{ dm}}{r_f}$$
 (11)

Substituting (10);

$$v_0^2 + \frac{m_0^2 r_0 GM}{m^2} \left(\frac{1}{r^2} - \frac{1}{r_f^2} \right) - 2GM \left(\frac{1}{r} - \frac{1}{r_f} \right) = 0$$
 (12)

Since $r_0 \!\! > \!\! > r_f$, if $r \!\! > r_0$, $r^{\text{-}1}$ and $r^{\text{-}2}$ terms can be neglected. Hence,

$$r_f = \frac{GM}{v_0^2} \left[\left(1 + \frac{m_0^2 r_0 v_0^2}{GMm^2} \right)^{1/2} - 1 \right].$$
 (0.8 pts)

To show that $r>r_0$ change in the linear momentum of the ordinary star in its reference frame:

$$-\frac{GMm}{r^2} + mr\omega^2 - m\frac{dv_r}{dt} = -v_0 \frac{dm_{gas}}{dt}$$
 (13)

and (13) implies the existence of an outward force initially and hence r starts growing. Using (7) one can write

$$mr\omega^2 = \frac{m_0^2 r_0^4 \omega_0^2}{mr^3}$$
.

Hence,
$$\frac{Gravitational\ force}{Centrifugal\ force} \alpha \text{ m}^2 r$$
. (0.4 pts)

where m is definitely decreasing. If r starts decreasing at some time also, this ratio starts decreasing, which is a contradiction.

So
$$r > r_0$$
. (0.4 pts)